

Dispersing Fermi-Ulam Models

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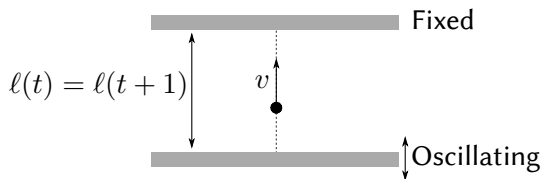
joint work with
Dmitry Dolgopyat · UMD

The Dynamical Systems, Ergodic Theory, and Probability
conference dedicated to the memory of Nikolai Chernov

UAB 20-05-2015

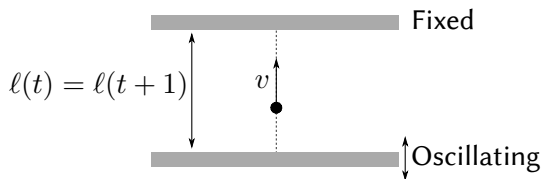
One-dimensional mechanical model

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Fermi-Ulam model

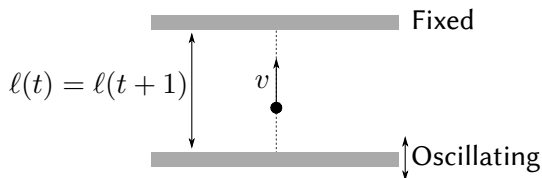
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Elastic collisions with the walls + free motion in between
Energy is **not** preserved at collisions

Fermi-Ulam model

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Goal: describe the long term energy distribution.

Conjecture (Fermi-Ulam¹⁹⁵¹)

“One expects that, after sufficiently long times, the average velocity of the point will become very large [...]. The tendency towards equipartition of energy would imply this.”

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Theorem (Pustynnikov¹⁹⁸³)

If $\ell(t)$ is analytic, all trajectories have bounded energy

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Theorem (Pustynnikov, Laederich-Levi, R. Douady¹⁹⁸³)

If $\ell(t) \in C^{4+\epsilon}$, all trajectories have bounded energy

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KAM-type result (main ingredient: Moser's small twist Theorem)

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Theorem (Zharnitsky¹⁹⁹⁹)

There exist $\ell(t) \in C^0$ so that some trajectories have unbounded energy

The collision map

Record data at every collision with the oscillating wall:

time of collision t

post-collisional velocity v

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time of collision	$t \in \mathbb{R}/\mathbb{Z}$
post-collisional velocity	$v \geq -\dot{\ell}(t)$

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$$M = \mathbb{T} \times \mathbb{R}^+ \quad f : (t_n, w_n) \mapsto (t_{n+1}, w_{n+1})$$

f is an exact twist map, $f_*\omega = \omega := w \, dt \wedge dw$ [note $\int_M \omega = \infty$]

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$$f(t_n, w_n) = (t_n + \delta t(t_n, w_n), w + \dot{\ell}(t_n) - \dot{\ell}(t_{n+1})),$$

where for $w > \|\dot{\ell}\|$:

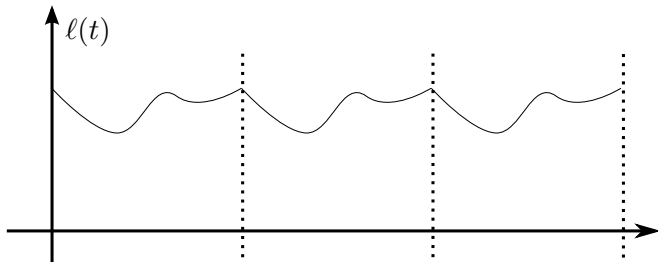
$$\delta t(t, w) = \frac{\ell(t) + \ell(t + \delta t(t, w))}{w - \dot{\ell}(t)} = \mathcal{O}(w^{-1})$$

for $w \rightarrow \infty$, f is close to integrable $(t, w) \mapsto (t + w^{-1}, w)$

Our general standing assumptions

Recall: $\ell(t)$ is distance between the two walls

$\ell(t) \in C^5[0, 1]$ strictly positive, $\ell(0) = \ell(1)$, $\dot{\ell}(0) \neq \dot{\ell}(1)$.



Early results: asymptotic normal form

Define *fundamental parameter* [\sim discontinuity]:

$$\Delta_\ell = \left[\ell(0) \int_0^1 \ell(s)^{-2} ds \right] \cdot (\dot{\ell}(0) - \dot{\ell}(1))$$

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Theorem (—Dolgopyat²⁰¹²)

There exists $R \simeq [0, 1] \times \mathbb{R}^+$, $R \subset M$ and coordinates (θ, I) on R with $I \sim v$ such that the first return map F of f on R is a $\mathcal{O}_5(I^{-1})$ -perturbation of $\hat{F} : (\theta, I) \mapsto (\bar{\theta}, \bar{I})$ where

$$\bar{\theta} = \theta - I \pmod{1} \qquad \bar{I} = I + \Delta_\ell(\bar{\theta} - 1/2)$$

The piecewise linear Standard Map I

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\hat{F} is a piecewise affine map: $\text{Tr } d\hat{F} \equiv 2 - \Delta_\ell$.

Corollary (of the Normal Form)

If $\Delta_\ell \in (0, 4)$ the dynamics of the FUM is asymptotically elliptic; the dynamics can be described [up to $\mathcal{O}(I^{-1})$] as a piecewise isometry.

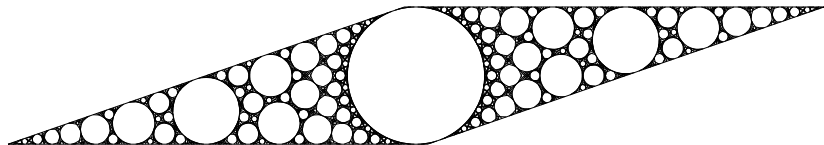
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If $\Delta_\ell \notin [0, 4]$ the dynamics of the FUM is asymptotically hyperbolic (i.e. \exists invariant cone fields for sufficiently large energies)

Moreover, let $\mathcal{E} = \{(t_0, w_0) \text{ s.t. } \lim_{n \rightarrow \infty} w_n = \infty\}$, then:

$$\text{mes } \mathcal{E} = 0 \qquad \text{HD } \mathcal{E} = 2,$$

recurrence, *stopped* CLT...

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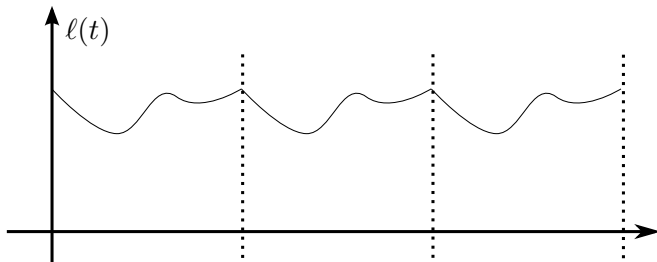
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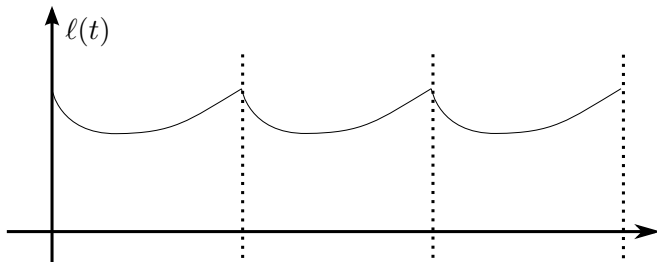
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Recall: $\ell(t)$ is distance between the two walls
 $\ell(t) \in C^5[0, 1]$, $\ell > 0$, $\ell(0) = \ell(1)$, $\dot{\ell}(0) \neq \dot{\ell}(1)$
+ Convexity assumption: $\ddot{\ell} \geq \kappa > 0$



Dispersing Fermi-Ulam models

Definition

A FUM is *dispersing* if $\ddot{\ell} \geq \kappa > 0$ (in particular $\Delta_\ell < 0$)

Theorem (—Dolgopyat²⁰¹⁵)

There exists a discrete set E s.t. if $\Delta_\ell \notin E$, the dynamics f of a dispersing FUM is ergodic.

e.g. the energy of a.e. trajectory can infinitely often be arbitrarily large and arbitrarily small.

Hints of the proof: hyperbolicity

Lemma

The dynamics f of a dispersing FUM is cone-hyperbolic, i.e

$$\exists \mathcal{K}^{u,s}(x) \subset T_x M \text{ s.t. } Df\mathcal{K}^u(x) \subset \mathcal{K}_{f(x)}^u, Df^{-1}\mathcal{K}^s(x) \subset \mathcal{K}_{f^{-1}(x)}^s$$

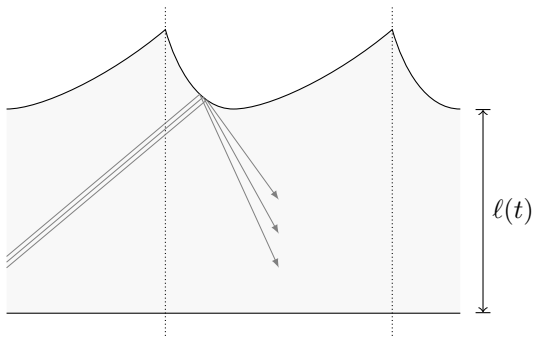
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Proof: f preserves the positive cone in (suitable) Jacobi coords \square



Hints of the proof: hyperbolicity

However:

$$\min_{v \in \mathcal{K}^u(x)} \frac{\|Df v\|_p}{\|v\|_p} \geq 1 + 2\tau(x)\kappa/w$$

where $\tau(x)$ is the flight time before the next collision and $\|\cdot\|_p$ is the p -metric

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- Large velocities
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Lemma

The induced map $\hat{f} : \hat{M} \rightarrow \hat{M}$ is uniformly hyperbolic.

Hints of the proof: from hyperbolicity to ergodicity

Ergodicity = local ergodicity (via Hopf argument)
+ combinatorial information

Ingredients

- Growth Lemma
- Distortion bounds for the (un)stable manifolds
- Absolute continuity of (un)stable manifolds

Our system has **infinite volume!**

Hints of the proof: the Growth Lemma

Unstable manifolds are *expanded* by hyperbolicity, but *fragmented* by singularities



\mathcal{W}^u



$f\mathcal{W}^u$



$f^2\mathcal{W}^u$

Growth Lemma

Let $l_n(x) = |\text{c.c. of } f^n\mathcal{W} \ni f^n x|$: if $n > C|\log|\mathcal{W}||$,

$$\mathbb{P}_{\mathcal{W}}(l_n(x) < \varepsilon) < C\varepsilon$$

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Theorem (—P.I. Tóth²⁰¹³)

Growth Lemma for planar billiards with corner points, finite horizon

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Theorem (—P.I. Tóth^{2013,WIP})

Growth Lemma for planar billiards with corner points, infinite horizon

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Hints of the proof: the Growth Lemma II

Fact

Subexponential complexity bounds on compact sets and at infinity (i.e. for the normal form) does not necessarily imply the Growth Lemma

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Additional assumption

Bounded complexity for the normal form (holds for $\Delta_\ell \notin E$)

\Rightarrow Growth Lemma

Thank you!